Math 216 — Second Midterm 20 March, 2017

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- 1. [14 points] Find real-valued solutions for each of the following, as indicated. (Note that minimal partial credit will be given on this problem.)
 - **a.** [7 points] Solve $\frac{1}{3}y'' + 2y' + 3y = 2t$, y(0) = 0, $y'(0) = \frac{4}{3}$.

Solution: The algebra may be easier if we first multiply by 3, obtaining y''+6y'+9y=6t. The characteristic equation for this is $\lambda^2 + 6\lambda + 9 = (\lambda + 3)^2 = 0$, so $\lambda = -3$ twice, and the homogeneous solution is $y_c = c_1 e^{-3t} + c_2 t e^{-3t}$. To find y_p we use the method of undetermined coefficients, guessing $y_p = At + B$. Then, plugging in,

$$6A + 9At + 9B = 6t,$$

so that $A = \frac{2}{3}$ and $B = -\frac{4}{9}$. The general solution is

$$y = c_1 e^{-3t} + c_2 t e^{-3t} + \frac{2}{3}t - \frac{4}{9}.$$

Applying the initial conditions, we have $y(0) = c_1 - \frac{4}{9} = 0$, so $c_1 = \frac{4}{9}$, and $y'(0) = -3c_1 + c_2 + \frac{2}{3} = c_2 - \frac{2}{3} = \frac{4}{3}$, so that $c_2 = 2$. Thus

$$y = \frac{4}{9}e^{-3t} + 2te^{-3t} + \frac{2}{3}t - \frac{4}{9}.$$

b. [7 points] Find the general solution to $y'' + 2y' + 5y = 2te^{-t}$.

Solution: The general solution will be $y = y_c + y_p$, where y_c solves the complementary homogenous problem and y_p is a particular solutions. For y_c we guess $y = e^{\lambda t}$, so that $\lambda^2 + 2\lambda + 5 = (\lambda + 1)^2 + 4 = 0$, and $\lambda = -1 \pm 2i$. Thus $y_c = c_1 e^{-t} \cos(2t) + c_2 e^{-t} \sin(2t)$. For y_p we use the method of undetermined coefficients, taking $y_p = (At+B)e^{-t}$. Plugging in, we have

$$(At - 2A + B)e^{-t} + 2(-At + A - B)e^{-t} + 5(At + B)e^{-t} = 3te^{-t}.$$

Collecting terms in e^{-t} and te^{-t} , we have 4B = 0 and 4A = 2. Thus B = 0 and $A = \frac{1}{2}$, and

$$y = c_1 e^{-t} \cos(2t) + c_2 e^{-t} \sin(2t) + \frac{1}{2} t e^{-t}$$

- **2**. [14 points] Find each of the following, providing an explicit formula where appropriate. (*Note that minimal partial credit will be given on this problem.*)
 - **a.** [5 points] $Y(s) = \mathcal{L}\{y(t)\}$ if $y'' + 4y' + 20y = 3\sin(2t), y(0) = 1, y'(0) = 2.$

Solution: Transforming both sides of the equation, we have

$$s^{2}Y - s - 2 + 4(sY - 1) + 20Y = \frac{6}{s^{2} + 4}$$

so that

$$Y = \frac{s+6}{s^2+4s+20} + \frac{6}{(s^2+4)(s^2+4s+20)}.$$

- b. [5 points] $\mathcal{L}^{-1}\left\{\frac{s}{s^2+4s+5}\right\}$ Solution: This is $\mathcal{L}^{-1}\left\{\frac{s}{s^2+4s+5}\right\} = \mathcal{L}^{-1}\left\{\frac{(s+2)-2}{(s+2)^2+1}\right\}$ $= \mathcal{L}^{-1}\left\{\frac{s+2}{(s+2)^2+1}\right\} - \mathcal{L}^{-1}\left\{\frac{2}{(s+2)^2+1}\right\}$ $= e^{-2t}\cos(t) - 2e^{-2t}\sin(t).$
 - c. [4 points] Using the integral definition of the Laplace transform, derive the transform rule $\mathcal{L}\{u_c(t)f(t-c)\} = e^{-sc}F(s)$ for a function f(t) with transform $L\{f(t)\} = F(s)$. (Recall $u_c(t)$ is the unit step function at t = c, $u_c(t) = \begin{cases} 0, & 0 < t < c \\ 1, & t \ge c \end{cases}$.)

Solution: The integral definition is $\mathcal{L}\{u_c(t)f(t-c)\} = \int_0^\infty e^{-st}u_c(t)f(t-c) dt$. Noting that $u_c(t)$ is zero for t < c, we may rewrite this as an integral with lower bound t = c. With the substitution w = t - c, we have

$$\int_0^\infty e^{-st} u_c(t) f(t-c) \, dt = \int_c^\infty e^{-st} f(t-c) \, dt = \int_0^\infty e^{-s(w+c)} f(w) \, dw$$
$$= e^{-cs} \int_0^\infty e^{-sw} f(w) \, dw = e^{-cs} F(s).$$

- **3**. [14 points] Use Laplace transforms to solve each of the following.
 - **a**. [7 points] $y'' + 4y' + 4y = 2e^{-2t}, y(0) = 1, y'(0) = 0.$

Solution: Taking the Laplace transform of both sides of the equation, we have with $Y = \mathcal{L}\{y\},\$

$$s^{2}Y - s + 4sY - 4 + 4Y = \frac{2}{s+2}$$
, or $Y = \frac{2}{(s+2)^{3}} + \frac{s+4}{(s+2)^{2}}$.

From transforms 3 and C from the table, we see that the first term in Y will invert as $\mathcal{L}^{-1}\left\{\frac{2}{(s+2)^3}\right\} = t^2 e^{-2t}$. To do the second, we use partial fractions: $\frac{s+4}{(s+2)^2} = \frac{A}{s+2} + \frac{B}{(s+2)^2}$. Clearing the denominators, we have s + 4 = A(s+2) + B, so that with s = -2 we find B = 2. Then s = 0 requires that A = 1, so that $\mathcal{L}^{-1}\left\{\frac{s+4}{(s+2)^2}\right\} = \mathcal{L}^{-1}\left\{\frac{1}{s+2} + \frac{2}{(s+2)^2}\right\} = e^{-2t} + 2te^{-2t}$. Combining this with the first result, we have

$$y = \mathcal{L}^{-1}\{Y\} = t^2 e^{-2t} + e^{-2t} + 2t e^{-2t}$$

b. [7 points]
$$y'' + 3y' = \begin{cases} 12, & 0 \le t < 2\\ 0, & t \ge 2 \end{cases}, y(0) = 0, y'(0) = 0.$$

Solution: We want to transform both sides of the equation; the right-hand side we can do by using the definition of the transform, or by noting that the differential equation may be written as $y'' + 3y' = 12 - 12u_2(t)$. Transforming the equation using transform 6 in the table, we have

$$s^{2}Y + 3sY = \frac{12}{s} - \frac{12e^{-2s}}{s}$$
, so that $Y = \frac{12}{s^{2}(s+3)} - \frac{12e^{-2s}}{s^{2}(s+3)}$

To find y we need to invert the transform of $\frac{12}{s^2(s+3)}$. We decompose this with partial fractions: $\frac{12}{s^2(s+3)} = \frac{A}{s} + \frac{B}{s^2} + \frac{C}{s+3}$. Clearing the denominators, $12 = As(s+3) + B(s+3) + Cs^2$. If s = 0, B = 4; if $s = -3, C = \frac{4}{3}$. Then, if $s = -2, 12 = -2A + 4 + \frac{16}{3}$, so that $A = -\frac{4}{3}$. Thus $\mathcal{L}^{-1}\{\frac{12}{s^2(s+3)}\} = \mathcal{L}^{-1}\{-\frac{4}{3s} + \frac{4}{s^2} + \frac{4}{3(s+3)}\} = -\frac{4}{3} + 4t + \frac{4}{3}e^{-3t}$, and, using this and transform 6 from the table, we have

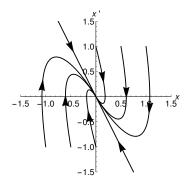
$$y = \mathcal{L}^{-1}\{Y\} = -\frac{4}{3} + 4t + \frac{4}{3}e^{-3t} - \left(-\frac{4}{3} + 4(t-2) + \frac{4}{3}e^{-3(t-2)}\right)u_2(t).$$

4. [14 points] Consider a mass-spring system modeled by

$$x'' + 4x' + \alpha x = 0$$

a. [5 points] Suppose that the phase portrait for the system is that shown to the right, below. For what values of α , if any, will the system have this type of behavior? Explain.

Solution: The characteristic equation of the equation is $\lambda^2 + 4\lambda + \alpha = (\lambda + 2)^2 + \alpha - 4 = 0$. The behavior shown in the phase portrait is that of a critically damped system, with a repeated eigenvalue and single eigenvector. This occurs when $\alpha = 4$. Thus we conclude that $\alpha = 4$.

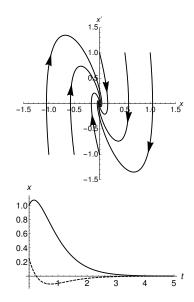


b. [3 points] For what values of α , if any, will the system be underdamped? Critically damped? Overdamped? Explain how you obtain your answers.

Solution: The work from (a) shows that for $\alpha = 4$ the system is critically damped. It is underdamped when it has a complex conjugate pair of roots, which will occur when $\alpha > 4$. It is overdamped when $\alpha < 4$.

c. [6 points] Let $\alpha = 6$. How will the phase portrait for the system in this case differ from that given in (a)? Sketch the phase portrait for this case. In a separate graph, sketch representative solutions x(t) as functions of time for the case $\alpha = 4$. (Note that you do not need to solve the problem to do this.)

Solution: From the work above, when $\alpha = 6$ the system will be underdamped. We can also see from the critically damped case in part (a) that the critical point (0,0) will be a sink and that trajectories will spiral in to the origin in a clockwise direction. This is shown in the figure to the right. When $\alpha = 4$ and we have the phase portrait shown in (a), the solution curves must converge to zero, with the possibility of crossing the *t*-axis (but not repeatedly), as shown in the lower figure.



- **5.** [15 points] For each of the following, identify the statement as true or false by circling "True" or "False" as appropriate, and provide a short (one or two sentence) explanation indicating why that answer is correct.
 - **a**. [3 points] For the system $x' = -xy + y^2$, $y' = x^2 xy$, the nonlinear trajectory in the phase plane with x(0) = -3 and y(0) = 0 lies on a circle centered on the origin.

True False

Solution: The system gives $\frac{dy}{dx} = \frac{y'}{x'} = \frac{x(x-y)}{-y(x-y)} = -\frac{x}{y}$. Separating and integrating, we get $y \, dy = -x \, dx$, so that $x^2 + y^2 = 2c$.

b. [3 points] For a linear differential operator $L = \frac{d^2}{dt^2} + p(t)\frac{d}{dt} + q(t)$, if y_1 and y_2 are different functions satisfying $L[y_1] = L[y_2] = g(t) \neq 0$, then, for any constants c_1 and c_2 , $y = c_1y_1 - c_2y_2$ satisfies L[y] = 0.

True False

Solution: Relying on the linearity of the operator,

$$L[y] = L[c_1y_1 - c_2y_2] = c_1L[y_1] - c_2L[y_2] = c_1g(t) - c_2g(t) = (c_1 - c_2)g(t)$$

which is zero only if $c_1 - c_2 = 0$.

c. [3 points] The solution to a differential equation my'' + ky = F(t) modeling the motion y of an undamped mechanical spring system with a periodic external force $F(t) = F_0 \cos(\omega t)$ can always be written as $y = A \cos(\omega_0 t - \delta_1) + B \cos(\omega t - \delta_2)$, a sum of two oscillatory terms. (A, B, ω_0 , δ_1 and δ_2 are constants.)

True False

True

Solution: This is only true if the forcing frequence ω is not equal to the natural frequency of the system, $\omega_0 = \sqrt{k/m}$. If $\omega = \omega_0$, we will have a growing solution $y = A \cos(\omega_0 t - \delta_1) + Bt \cos(\omega_0 - \delta_2)$.

d. [3 points] If $\lambda^2 + p\lambda + q = 0$ is the characteristic equation of a constant-coefficient linear differential equation L[y] = g(t), then solving for $Y(s) = \mathcal{L}\{y(t)\}$ will result in an expression involving a product of $(s^2 + ps + q)^{-1}$ with other terms.

False

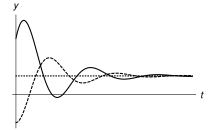
Solution: The characteristic equation tells us that the differential equation is L[y] = y'' + py' + qy = g(t). The transform of this is $s^2Y - sy(0) - y'(0) + psY - y(0) + qY = G(s)$, and solving for Y will give the indicated result.

e. [3 points] If $f(t) \neq 0$ has Laplace transform $\mathcal{L}{f(t)} = F(s)$ and $g(t) = \begin{cases} f(t), & 0 < t < c \\ 0, & t \ge c \end{cases}$, then $\mathcal{L}{g(t)} = (1 - e^{-sc})F(s)$. True

Solution: The easiest way to see this is to look for $h(t) = \mathcal{L}^{-1}\{(1 - e^{-sc})F(s)\} = f(t) - f(t-c)u_c(t)$. Then, by definition, $g(t) = (1 - u_c(t))f(t)$, and in general $h(t) \neq g(t)$.

- **6**. [14 points] In the following, we consider the behavior of solutions to a linear, second-order, constant-coefficient differential equation with a forcing term.
 - **a**. [5 points] Write a differential equation of this type that could have the three solution curves given to the right. Explain how you know your answer is correct.

Solution: The constant solution shows that we have a non-zero equilibrium solution, so the forcing term is g(t) = k, a constant. Then the two non-constant solutions show an oscillatory transient with decaying amplitude, so the characteristic equation of the



differential equation must have complex roots with a negative real part. Thus any equation of the form y'' + ay' + by = k, where a > 0 and $a^2 - 4b < 0$ (and, because the equilibrium is positive, k > 0) will produce the desired result. One such is y'' + y' + y = 1.

b. [6 points] Now suppose that the general solution to the problem is $y = (c_1 + c_2 t + t \ln(t))e^{-t}$. What is the differential equation, including the forcing term? Explain.

Solution: Because the problem is linear we know that the general solution has the form $y = y_c + y_p$, where y_c is the solution to the complementary homogeneous problem and y_p is a solution to the problem with forcing. Because it is constant-coefficient and second-order, the solution y_c has terms of the form $e^{\lambda t}$ or $te^{\lambda t}$ (where λ may be zero or complex), so the homogeneous solution here must be $y_c = c_1 e^{-t} + c_2 t e^{-t}$. This requires that $\lambda = -1$, twice, so the characteristic equation is $\lambda^2 + 2\lambda + 1 = 0$, and the linear differential operator is L[y] = y'' + 2y' + y. Then $y_p = t \ln(t)e^{-t}$. We can find g(t) by plugging this into L[y]; to do this, we calculate $y'_p = -t \ln(t)e^{-t} + (\ln(t) + 1)e^{-t}$ and $y''_p = -(y_p)' + \frac{1}{t}e^{-t} - (\ln(t) + 1)e^{-t}$. Then

$$L[y_p] = y_p'' + 2y_p' + y_p$$

= $(\frac{1}{t}e^{-t} - (\ln(t) + 1)e^{-t}) + (-t\ln(t)e^{-t} + (\ln(t) + 1)e^{-t}) + t\ln(t)e^{-t}$
= $t^{-1}e^{-t} = g(t).$

Thus the equation and forcing are $y'' + 2y' + y = t^{-1}e^{-t}$.

c. [3 points] If you were finding, by hand, the general solution given in (b), what method or methods could you use? In these methods, what form do you guess for the solution?

Solution: Because the problem is linear, we know that we will be finding the solution to the complementary homogeneous problem and then finding a particular solution. Because it is constant-coefficient, the former will always be done by finding the solution to the eigenvalue problem obtained by looking for solutions of the form $y = e^{\lambda t}$ (or, $\mathbf{x} = \mathbf{v}e^{\lambda t}$ for the equivalent system of two first-order equations). The forcing term (and particular solution) do not admit use of the method of undertermined coefficients, so we would use variation of parameters to guess $y_p = u_1(t)e^{-t} + u_2(t)te^{-t}$.

7. [15 points] In lab 3 we considered a nonlinear system modeling a laser with a slightly varying gain rate, which we rewrite slightly in this problem as

$$N' = \gamma(A - N(1 + P))$$
$$P' = P(N - 1)$$

with $A = A_0 + \epsilon \cos(\omega t)$.

a. [5 points] If A is constant, the system has a critical point (N, P) = (1, A - 1). Let N = 1 + u, $P = A_0 - 1 + v$, and $A = A_0 + \epsilon \cos(\omega t)$ and find a linear system in u and v by assuming that u, v and ϵ are all very small.

Solution: Substituting these values into the system, we have

$$u' = \gamma(A_0 + \epsilon \cos(\omega t) - (1 + u)(A_0 + v))$$
$$= \gamma(-A_0 u - v) + \epsilon \gamma \cos(\omega t) - \gamma uv$$
$$v' = (A_0 - 1 + v)u = (A_0 - 1)u + uv$$

Discarding the (very, very,) very small terms uv, we have

$$u' = \gamma(-A_0u - v) + \epsilon\gamma\cos(\omega t), \quad v' = (A_0 - 1)u$$

b. [5 points] The system you obtained in (a) can be rewritten, for some constants α and β , as $v'' + \alpha v' + \beta v = \epsilon \beta \cos(\omega t)$. Find the steady-state response to this rewritten form.

Solution: The steady-state response is v_p . Letting $v_p = B\cos(\omega t) + C\sin(\omega t)$ and plugging into the equation, we have

$$-B\omega^2\cos(\omega t) - C\omega^2\sin(\omega t) -B\alpha\omega\sin(\omega t) + C\alpha\omega\cos(\omega t) + B\beta\cos(\omega t) + C\beta\sin(\omega t) = \epsilon\beta\cos(\omega t).$$

Collecting terms in $\cos(\omega t)$ and $\sin(\omega t)$, we have two equations for B and C,

$$(-\omega^{2} + \beta)B + \alpha\omega C = \epsilon\beta$$
$$-\alpha\omega B + (-\omega^{2} + \beta)C = 0.$$

Multiplying the first by $\alpha\omega$ and the second by $-\omega^2 + \beta$ and adding, we find

$$C = \frac{\epsilon\beta\alpha\omega}{(-\omega^2 + \beta)^2 + (\alpha\omega)^2}, \quad B = \frac{\epsilon\beta(-\omega^2 + \beta)}{(-\omega^2 + \beta)^2 + (\alpha\omega)^2},$$

where the result for B follows by solving the second equation for C in terms of B. Then the steady-state response is $v = B\cos(\omega t) + C\sin(\omega t)$ for these B and C.

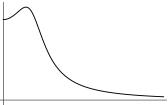
Problem 7, continued.

c. [5 points] Suppose that the steady-state solution that you obtained in (b) was, for some constant b with |b| < 1, $v_{ss} = \frac{b\omega}{(1-\omega^2)^2 + (b\omega)^2} \cos(\omega t) + \frac{1-\omega^2}{(1-\omega^2)^2 + (b\omega)^2} \sin(\omega t)$. Find the amplitude of the oscillation and explain why the solution exhibits resonance behavior.

Solution: Note that this is of the form $v_{ss} = B\cos(\omega t) + C\sin(\omega t) = R\cos(\omega t - \delta)$, with $R = \sqrt{B^2 + C^2}$. Thus the amplitude of the response is

$$\begin{split} R &= \sqrt{\left(\frac{b\omega}{(1-\omega^2)^2 + (b\,\omega)^2}\right)^2 + \left(\frac{1-\omega^2}{(1-\omega^2)^2 + (b\,\omega)^2}\right)^2} \\ &= \sqrt{\frac{(1-\omega)^2 + (b\omega)^2}{((1-\omega^2)^2 + (b\,\omega)^2)^2}} = \frac{1}{\sqrt{(1-\omega^2)^2 + (b\,\omega)^2}}. \end{split}$$

We expect this to be a function similar to that shown below



(This can be deduced from the form of R as well: if $\omega = 0$, R = 1; if $\omega = 1$, R = 1/|b| > 1, and as $\omega \to \infty$, $R \to 0$.) Thus, for some intermediate value of ω (a bit less than one) there is a maximum response amplitude, which is what we mean by resonance.

Formulas, Possibly Useful

- Some Taylor series, taken about x = 0: $e^x = \sum_{n=0}^{\infty} \frac{x^n}{n!}$; $\cos(x) = \sum_{n=0}^{\infty} (-1)^n \frac{x^{2n}}{(2n)!}$; $\sin(x) = \sum_{n=0}^{\infty} (-1)^n \frac{x^{2n+1}}{(2n+1)!}$. The series for $\ln(x)$, taken about x = 1: $\ln(x) = \sum_{n=1}^{\infty} (-1)^{n+1} \frac{(x-1)^n}{n}$.
- Some integration formulas: $\int u v' dt = u v \int u' v dt$; thus $\int t e^t dt = t e^t e^t + C$, $\int t \cos(t) dt = t \sin(t) + \cos(t) + C$, and $\int t \sin(t) dt = -t \cos(t) + \sin(t) + C$.

	f(t)	F(s)
1.	1	$\frac{1}{s}, s > 0$
2.	e^{at}	$\frac{1}{s-a}, s > a$
3.	t^n	$\frac{n!}{s^{n+1}}$
4.	$\sin(at)$	$\frac{a}{s^2 + a^2}$
5.	$\cos(at)$	$\frac{s}{s^2 + a^2}$
6.	$u_c(t)$	$\frac{e^{-cs}}{s}$
7.	$\delta(t-c)$	e^{-cs}
А.	f'(t)	s F(s) - f(0)
A.1	f''(t)	$s^2 F(s) - s f(0) - f'(0)$
A.2	$f^{(n)}(t)$	$s^n F(s) - \dots - f^{(n-1)}(0)$
В.	$t^n f(t)$	$(-1)^n F^{(n)}(s)$
С.	$e^{ct}f(t)$	F(s-c)
D.	$u_c(t) f(t-c)$	$e^{-cs} F(s)$
E.	f(t) (periodic with period T)	$\frac{1}{1 - e^{-Ts}} \int_0^T e^{-st} f(t) dt$

Some Laplace Transforms