$\begin{array}{l} \text{Math 216} - \text{Second Midterm} \\ \text{ 19 March, 2018} \end{array}$

This sample exam is provided to serve as one component of your studying for this exam in this course. Please note that it is not guaranteed to cover the material that will appear on your exam, nor to be of the same length or difficulty. In particular, the sections in the text that were covered on this exam may be slightly different from those covered by your exam.

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- **1**. [15 points] Find explicit, real-valued solutions to each of the following, as indicated. For this problem, use Laplace transforms, **not** some other solution technique.
 - **a.** [8 points] $y'' + 2y' + y = e^{-t}, y(0) = 0, y'(0) = 1.$

Solution: Applying the forward transform, letting $Y = \mathcal{L}\{y\}$, we have $(s^2+2s+1)Y-1 = \frac{1}{s+1}$, so that

$$Y = \frac{1}{(s+1)^2} + \frac{1}{(s+1)^3}$$

Note that, using rules (2) and (B) from the table, we have $\mathcal{L}^{-1}\left\{\frac{1}{(s+1)^2}\right\} = \mathcal{L}^{-1}\left\{-\frac{d}{ds}\frac{1}{(s+1)}\right\} = te^{-t}$, and $\mathcal{L}^{-1}\left\{\frac{1}{(s+1)^3}\right\} = \mathcal{L}^{-1}\left\{-\frac{1}{2}\frac{d}{ds}\frac{1}{(s+1)^2}\right\} = \frac{1}{2}t^2e^{-t}$. Thus,

$$y = \mathcal{L}^{-1}\left\{\frac{1}{(s+1)^2}\right\} + \mathcal{L}^{-1}\left\{\frac{1}{(s+1)^3}\right\}$$
$$= te^{-t} + \frac{1}{2}t^2e^{-t}.$$

b. [7 points] y'' + 6y' + 13y = 0, y(0) = 1, y'(0) = 0.

Solution: Applying the forward transform, letting $Y = \mathcal{L}\{y\}$, we have $(s^2 + 6s + 13)Y - s - 6 = 0$, so that

$$Y = \frac{s+3}{(s+3)^2+4} + \frac{3}{(s+3)^2+4}$$

We can invert both of these with rules (4), (5), and (C). We have

$$y = \mathcal{L}^{-1}\{Y\} = \mathcal{L}^{-1}\left\{\frac{s+3}{(s+3)^2+4} + \frac{3}{(s+3)^2+4}\right\}$$
$$= e^{-3t}\cos(2t) + \frac{3}{2}e^{-3t}\sin(2t).$$

2. [14 points] Fill in the missing portions of each of the following transforms. Briefly explain how you obtain your work.

a. [7 points]
$$\mathcal{L}\left\{\begin{cases} 0, & 0 \le t < 1\\ 1, & 1 \le t < 5 \end{cases}\right\} = \frac{1}{s}(e^{-s} - e^{-5s}) + \underline{e^{-5s} \frac{1}{s+1}} \\ e^{-(t-5)}, & t \ge 5 \end{cases}$$

Solution: Letting f be the indicated function and breaking the integral on the piecewise definitions, we have

$$\mathcal{L}\{f\} = \int_{1}^{5} e^{-st} dt + \int_{5}^{\infty} e^{-(t-5)} e^{-st} dt.$$

The first of these integrals gives the terms provided. The second is

$$\int_{5}^{\infty} e^{5} e^{-(s+1)t} dt = \lim_{b \to \infty} \left(-\frac{e^{5}}{s+1}\right) e^{-(s+1)t} \Big|_{t=5}^{t=b} = \frac{e^{5}}{s+1} e^{-5s-5} = \frac{e^{-5s}}{s+1} e^{-5s-$$

b. [7 points]
$$\mathcal{L}^{-1}\left\{\frac{1}{s(s+1)(s^2+1)}\right\} = 1 - \frac{1}{2}\cos(t) + \frac{-\frac{1}{2}\sin(t) - \frac{1}{2}e^{-t}}{-\frac{1}{2}\sin(t) - \frac{1}{2}e^{-t}}$$

Solution: We can rewrite this with partial fractions:

$$\frac{1}{s(s+1)(s^2+1)} = \frac{A}{s} + \frac{B}{s+1} + \frac{Cs+D}{s^2+1}.$$

The inverse transforms of these will be, in order, A, Be^{-t} , $C\cos(t)$, and $D\sin(t)$. Thus from the provided partial answer we know that A = 1 and $C = -\frac{1}{2}$. Then, clearing the denominators and using these values, we have

$$1 = (s+1)(s^{2}+1) + Bs(s^{2}+1) + (-\frac{1}{2}s + D)s(s+1).$$

If s = -1, we have 1 = -2B, so that $B = -\frac{1}{2}$. If s = 1, we have $1 = 4 - 1 + (-\frac{1}{2} + D)(2) = 2 + 2D$. Thus $D = -\frac{1}{2}$, and we have the indicated result.

- **3**. [14 points] Find each of the following, as indicated.
 - **a.** [7 points] If a function f(t) has the Laplace transform $F(s) = \mathcal{L}\{f(t)\}$, use the integral definition of the Laplace transform to find the transform $\mathcal{L}\{\int_0^t f(x) dx\}$ in terms of F(s). (You may assume that $\int_0^\infty f(x) dx = L$, a finite value.)

Solution: We have, integrating by parts with $u = \int_0^t f(x) dx$ (so that u' = f(t)) and $v' = e^{-st}$ (so that $v = -\frac{1}{s}e^{-st}$), $\mathcal{L}\{\int_0^T f(t) dt\} = \int_0^\infty (\int_0^t f(x) dx) e^{-st} dt$ $= -\frac{1}{s} \lim_{b \to \infty} (\int_0^t f(x) dx) e^{-st} \Big|_{t=0}^{t=b} + \frac{1}{s} \int_0^\infty f(t)e^{-st} dt$ $= -\frac{1}{s} \int_0^0 f(x) dx + \lim_{b \to \infty} \frac{1}{s} L e^{-sb} + \frac{1}{s} F(s).$

The first term on the right-hand side is an integral over an interval of zero length, and so is zero, and in the limit as $b \to \infty$, the second vanishes because of the negative exponential e^{-sb} . Thus we have $\mathcal{L}\{\int_0^t f(x) dx\} = \frac{1}{s} F(s)$.

b. [7 points] Find an explicit expression for $Y = \mathcal{L}\{y\}$ if $y''' + 3y = t^2 e^{-4t} - e^{-2t} \cos(5t)$. (Note that you are not asked to solve the differential equation.)

Solution: Because of the linearity of the transform, we can calculate the transform of each term separately. We have $\mathcal{L}\{y'''\} = s^3Y - s^2y(0) - sy'(0) - y''(0)$ and $\mathcal{L}\{3y\} = 3Y$. To find $\mathcal{L}\{t^2e^{-4t}\}$, we note that $\mathcal{L}\{e^{-4t}\} = \frac{1}{s+4}$, so that, by rule (B), $\mathcal{L}\{t^2e^{-4t}\} = \frac{d^2}{ds^2}\frac{1}{s+4} = \frac{2}{(s+4)^3}$. To find $\mathcal{L}\{e^{-2t}\cos(5t)\}$, we use rule (C) and the transform $\mathcal{L}\{\cos(5t)\} = \frac{s}{s^2+25}$ to get $\mathcal{L}\{e^{-2t}\cos(5t)\} = \frac{s+2}{(s+2)^2+25}$. Putting these all together and solving for Y, we have

$$Y = \frac{s^2 y(0) + sy'(0) + y''(0)}{s^3 + 3} + \frac{2}{(s+4)^3(s^3 + 3)} - \frac{s+2}{((s+2)^2 + 25)(s^3 + 3)}$$

- **4**. [15 points] Find explicit, real-valued solutions for each of the following, as indicated. Do **not** use Laplace transform techniques on this problem.
 - **a**. [8 points] Find the general solution to $y'' + 2y' + 4y = e^{-t} + t^2$.

Solution: We first look for the general solution in the form $y = y_c + y_p$. For y_c , the guess $y = e^{rt}$ gives $r^2 + 2r + 4 = (r+1)^2 + 3 = 0$, so that $r = -1 \pm \sqrt{3}i$, and $y_c = c_1 e^{-t} \cos(\sqrt{3}t) + c_2 e^{-t} \sin(\sqrt{3}t)$.

To find y_p , we use the method of undetermined coefficients, looking for each of the forms on the right-hand side separately. For the e^{-t} term, we guess $y = ae^{-t}$, so that (1-2+4)a = 1, and $a = \frac{1}{3}$. For the t^2 term, we guess $y = a_0 + a_1t + a_2t^2$, so that $2a_2 + 2(2a_2t + a_1) + 4(a_2t^2 + a_1t + a_0) = t^2$. Collecting powers of t, we have $a_2 = \frac{1}{4}$, $a_1 = -\frac{1}{4}$, and $a_0 = 0$.

The general solution to the problem is therefore

$$y = c_1 e^{-t} \cos(\sqrt{3}t) + c_2 e^{-t} \sin(\sqrt{3}t) + \frac{1}{3}e^{-t} + \frac{1}{4}t^2 - \frac{1}{4}t.$$

b. [7 points] Solve $y'' + 5y' + 4y = 3\cos(2t), y(0) = 0, y'(0) = 1$

Solution: We look for the general solution as $y = y_c + y_p$. For y_c we have $y = e^{rt}$, and $r^2 + 5r + 4 = (r+4)(r+1) = 0$, so that $y_c = c_1 e^{-4t} + c_2 e^{-t}$.

For y_p we use the method of undetermined coefficients guess $y_p = a \cos(2t) + b \sin(2t)$. Plugging in, we have

$$-4a\cos(2t) - 4b\sin(2t) - 10a\sin(2t) + 10b\cos(2t) + 4a\cos(2t) + 4b\sin(2t) = 3\cos(2t).$$

Collecting terms in $\cos(2t)$ and $\sin(2t)$, we have 10b = 3, and -10a = 0; thus a = 0 and $b = \frac{3}{10}$. Our general solution is therefore

$$y = c_1 e^{-4t} + c_2 e^{-t} + \frac{3}{10} \sin(2t).$$

Applying the initial conditions, we have $y(0) = c_1 + c_2 = 0$, and $y'(0) = -4c_1 - c_2 + \frac{3}{5} = 1$. Substituting the first into the second, we have $3c_2 = \frac{2}{5}$, so that $c_2 = \frac{2}{15}$ and $c_1 = -\frac{2}{15}$. Thus

$$y = -\frac{2}{15}e^{-4t} + \frac{2}{15}e^{-t} + \frac{3}{10}\sin(2t).$$

5. [14 points] Consider the operators $T[y] = yy'' + 2y^2y'$ and $U[y] = t^2y'' - ty' - 3y$. **a**. [9 points] Show that T is nonlinear while U is linear.

Solution: We note that

$$T[cy] = cy(cy'') - 2c^2y^2(cy') = c^2(yy'' - 2cy^2y') \neq cT[y]$$

(because $cT[y] = c(yy'' + 2y^2y')$). Therefore T is not linear. (Other possible arguments include the calculations $T[y_1 + y_2] \neq T[y_1 + y_2]$, and $T[c_1y_1 + c_2y_2] \neq c_1T[y_1] + c_2T[y_2]$.) However,

$$U[c_1y_1 + c_2y_2] = t^2(c_1y_1'' + c_2y_2'') - t(c_1y_1' + c_2y_2') - 3(c_1y_1 + c_2y_2)$$

= $c_1(t^2y_1'' - ty_1' - 3y_1) + c_2(t^2y_2'' - ty_s' - 3y_2)$
= $c_1U[y_1] + c_2U[y_2].$

Thus U is linear.

Alternately, we said that a linear operator may also be written in the form $L = D^2 + D^2$ p(t)D + q(t), so that $L[y] = D^2[y] + p(t)D[y] + q(t)y$. For T, we have

$$T[y] = yy'' + 2y^2y' = yD^2[y] + 2y^2D[y] = y(D^2[y] + 2yD[y]),$$

which is not in the form of a linear operator, while

$$U[y] = t^{2}(D^{2}[y] - t^{-1}D[y] - 3y) = t^{2}(D^{2} - t^{-1}D - 3t^{-2})[y],$$

so that the operator $U = t^2 D^2 - tD - 3$. Note that this isn't in the form of L; it is an easy generalization from that, but isn't quite consistent with our understanding. It's best to make the argument above, or to consider the operator in the context of a differential equation: the equation U[y] = 0 can be written as L[y] = 0, with $L = D^2 - t^{-1}D - 3t^{-2}$, a linear operator.

b. [5 points] Show that $y_1 = t^{-1}$ and $y_2 = t^3$ constitute a fundamental set of solutions to the equation U[y] = 0. What is the general solution to U[y] = 0? (You may assume that t > 0.)

Solution: A fundamental solution set to a linear second-order homogeneous equation consists of two linearly independent solutions. Note that $U[y_1] = t^2(2t^{-3}) - t(-t^{-2}) - t(-t^{-2})$ $3t^{-1} = 3t^{-1} - 3t^{-1} = 0$, and $U[y_2] = t^2(6t) - t(3t^2) - 3t^3 = 6t^3 - 6t^3 = 0$, so y_1 and y_2 are solutions to U[y] = 0. Then

$$W[y_1, y_2] = \begin{vmatrix} t^{-1} & t^3 \\ -t^{-2} & 3t^2 \end{vmatrix} = 3t + t = 4t \neq 0,$$

so the solutions are linearly independent. Thus the general solution is $y = c_1 t^{-1} + c_2 t^3$.

- 6. [13 points] Consider the phase portrait shown to the right, which shows the phase portrait for a linear, second-order, constant coefficient, homogeneous differential equation L[y] = 0.
 - **a**. [7 points] Write a differential equation that could give this phase portrait. Explain how you obtain your solution, and why is it correct.

Solution: We note that the phase portrait shows a center, that is, trajectories are simple closed loops. This suggests that the solutions are sines and cosines, so that we should have an equation $y'' + \omega_0^2 y = 0$. Solutions to this are $y_1 = \cos(\omega_0 t)$ and $y_2 = \sin(\omega_0 t)$, so that the phase plane trajectories are given by $\mathbf{x}_1 = c_1 \begin{pmatrix} \cos(\omega_0 t) \\ -\omega_0 \sin(\omega_0 t) \end{pmatrix}$ and $\mathbf{x}_2 = c_2 \begin{pmatrix} \sin(\omega_0 t) \\ \omega_0 \cos(\omega_0 t) \end{pmatrix}$. We note that the vertical stretch of the shown trajectories appears to be twice that of the horizontal, so guess that $\omega_0 = 2$. Our equation is therefore y'' + 4y = 0.



b. [6 points] Suppose that we add a forcing term $f(t) = \cos(15t/8)$ to the equation, so that we are solving L[y] = f(t). Sketch an approximate solution curve with y(0) = 0, y'(0) = 1. Explain why your solution appears as it does.

Solution: Note that the forcing frequency $\omega = 15/8$ is close to the natural frequency of the system, $\omega_0 = 2$. So we expect to see at least a mild beats phenomenon. This is shown in the graph below.



7. [15 points] In our lab on lasers, we considered a linearization of the nonlinear model for the population inversion N and light intensity P. A critical point of the nonlinear system is (N, P) = (1, A - 1), and linearizing the system near this gives the linear system

$$u' = -\gamma(Au + v), \qquad v' = (A - 1)u,$$

where γ and A are constants.

a. [5 points] Rewrite this as a single, second-order equation in v.

Solution: Note that, from the second equation, $u = \frac{1}{A-1}v'$. Plugging this into the first equation, we have $\frac{1}{A-1}v'' + \frac{\gamma A}{A-1}v' + \gamma v = 0,$ or $v'' + \gamma Av' + \gamma (A-1)v = 0.$

b. [5 points] Suppose that for some α and β your equation from (a) is $v'' + \alpha v' + \beta v = 0$. Under what conditions on α and β will the solution for v be underdamped? Write down two real-valued linearly independent solutions to the equation in this case.

Solution: We note that solutions to this equation are $v = e^{rt}$ with $r = -\frac{\alpha}{2} \pm \frac{1}{2}\sqrt{\alpha^2 - 4\beta}$. This will give underdamping if $\alpha^2 - 4\beta < 0$, that is, if $\alpha^2 < 4\beta$. In terms of the constants we obtained in (a), this is $\gamma^2 A^2 < 4\gamma(A-1)$, or $\gamma < 4\frac{A-1}{A}^2$. The two solutions are $y_1 = e^{-\mu t} \cos(\nu t)$ and $y_2 = e^{-\mu t} \sin(\nu t)$, where $\mu = \frac{\alpha}{2} = \frac{\gamma A}{2}$ and $\nu = \frac{1}{2}\sqrt{4\beta - \alpha^2} = \frac{1}{2}\sqrt{4\gamma(A-1) - \gamma^2 A^2}$.

c. [5 points] Now suppose that we force the underdamped equation given in (b) with the periodic forcing term $f(t) = \cos(\omega t)$. Sketch a graph of the steady state solution of the problem. Explain why your graph has the form it does. If ω changes from very small to very large values, how would you expect your sketch to change? Explain.

Solution: The steady state solution will be the response to the forcing, because (as we see in (b)) the non-forced response decays to zero. Because y_1 and y_2 do not have the same form as f(t), we know that the steady state (particular) solution will have the form $v_p = a \cos(\omega t) + b \sin(\omega t) = R \cos(\omega t - \phi)$, so it will be a simple sinusoid:



If we vary ω , we expect that the frequency of the solution will change, and that its amplitude will also change. A reasonable guess is that the amplitude will initially increase, obtain a local maximum at an ω near $\nu = \frac{1}{2}\sqrt{4\beta - \alpha^2} = \frac{1}{2}\sqrt{4\gamma(A-1) - \gamma^2 A^2}$, and then decay to zero as ω becomes very large.